Seminarios CADEDIF U.C.M. Madrid, 30 Abril 2009 **ASYMPTOTIC BEHAVIOUR FOR** SMALL WIDTH OF INTERFACE **IN PHASE-FIELD MODEL** ANGELA JIMÉNEZ-CASAS Grupos Dinámica No Lineal y CADEDIF Universidad Pontificia Comillas de Madrid. Universidad Complutense de Madrid

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- $\tau = (\xi)^2$ thin-interface limit
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1. PHASE FIELD MODEL EQUA-TIONS

Stefan's Problem (solid-liquid)

• The evolution of the temperature , u(t,x), of the point $x \in \Omega \subset I\!\!R^N$ at time t

of a substance which may appear in two different phases.

• The evolution of the interphase Γ .

 $\Gamma(t) = \{ x \in \Omega \text{ such that } u(t, x) = 0 \},\$

The liquid phase is given by:

$$\Omega_1 = \{ x \in \Omega \text{ such that } u(t, x) > 0 \}$$

The solid phase is given by:

 $\Omega_2 = \{ x \in \Omega \text{ such that } u(t, x) < 0 \}.$

Enthalpy method o H-method:

balance heat is given by the diffusion equation

$$\frac{\partial}{\partial t}H(u) = k\Delta u \tag{0.1}$$

with k > 0, diffusivity constant and

$$H(u) = u + \frac{l}{2}\varphi$$
 enthalpy function.

where l > 0, latent heat

- φ is the known function, associated to the change phase This step function implies that we consider the linear interphase set Γ.
- But, the Stefan's model can not explain some phenomenons which appear in the equilibrium (supercooling), so we have to consider the interface set is not linear.
- If we consider a **plane** region of interface of width ξ
- We have a new unknown function φ, instead the step function of Stefan's model.

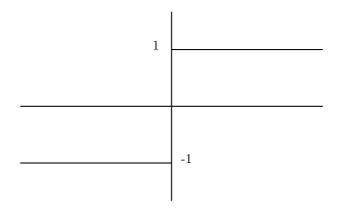


Figure 1: Step function of change phase,φ

♣ Interphase plane (non linear). Phase field function or Order parameter

• $\varphi(t, x)$ is the known function, associated to the change phase, Phase field function or Order parameter,

 φ is scalar function depends on time t and the position x and take different values in two different phases.

 $\varphi(t,x): I\!\!R^+ \times \Omega \longmapsto I\!\!R$

is local average of phase (solido - liquid).

From Landau-Ginzburg's theory, the free energy of sys-

tem is given by

$$F_u(\varphi) = \int [\frac{1}{2}\xi^2 (\nabla \varphi)^2 + \frac{1}{8}(\varphi^2 - 1)^2 - 2u\varphi]dx \qquad (0.2)$$

• The equilibrium equation. Euler-Lagrange

The system is in equilibrium if (u, φ) satisfies:

$$\begin{cases} 0 = \xi^2 \Delta \varphi + \frac{1}{2}(\varphi - \varphi^3) + 2u \\ 0 = \Delta u \end{cases}$$
(0.3)

together with the boundary conditions.

Phase field equations. Landau-Ginzburg

• The evolution of φ and u is given by the parabolic system

$$\begin{cases} \tau \varphi_t = w \Delta \varphi - f(\varphi) + 2u & \text{in } \Omega \times I\!\!R^+ \\ u_t + \frac{l}{2} \varphi_t = k \Delta u & \text{in } \Omega \times I\!\!R^+ \end{cases}$$
(0.4)

- + BOUNDARY CONDITIONS
- + INITIAL CONDITIONS

- $u(x,t) \equiv temperature of point \ x \in \Omega$ at time t.
- $\varphi(x,t) \equiv order \ parameter \ or \ phase \ field.$
- Ω is an open and bounded set in \mathbb{R}^N , $N \ge 1$, with regular boundary.
- $f(\varphi)$ is typically $\frac{1}{2}(\varphi^3 \varphi)$.
- l and $k \equiv$ are positive constants associated to latent heat (l) and thermal diffusivity (k).
- τ and $w \equiv$ are positive parameters related to time and length scales.
- $w = w(\xi)$ with ξ width of interface.

Phase Field in Biology/Industrial

• The phase-field can be seen as the density of bacterial collony or the mass of growing tumor.

Analogously, the diffusion field an stand for the density of nutrient. [13]

We show that this phase-field approach is suitable to describe homogeneous as well as heterogenous nucleation starting from several general hypotheses.

 \bullet Quantitative phase-field modeling of dendritic growth in

two and three dimensions [23]

• The phase-field can be see the dynamics of phase separation and coarsening of mixtures of three or more components.

In this case de function u(t, x) denote the concentration of the point x at time t, of one the components of the mixture.

2. ASYMPTOTIC BEHAVIOUR FOR SMALL WIDTH OF INTER-FACE ξ .

 $\clubsuit One-dimensional \ semilinear \ parabolic \ system \equiv \ Phase \\ Field \ model \\$

$$\begin{cases} \tau \varphi_t = w \varphi_{xx} - f(\varphi) + 2u, \ x \in (a, b) \\ u_t + \frac{l}{2} \varphi_t = k u_{xx}, & x \in (a, b) \\ \varphi'(a) = \varphi'(b) = 0 \\ u'(a) = u'(b) = 0 \\ \varphi(0, x) = w_0(x) \in H^1(a, b) \\ u(0, x) = u_0(x) \in L^2(a, b) \end{cases}$$
(0.5)

• A substance which may appear in different phases

- $f(\varphi) = \frac{1}{2}(\varphi^3 \varphi)$ (in two different phases)
- $u(t, x) \equiv$ temperature of the point x at time t
- $\varphi(t, x) \equiv$ is the phase field function or order parameter, (local phase average).
- $l \equiv latent head, k \equiv thermal diffusivity.$
- $\tau \equiv time \ scale$.
- $w \equiv length \ scale \ (\ w = w(\xi), \xi \equiv interface \ width).$
- G. Caginalp 1986, 1990 and 1991, P.C. Fife 1988 and 1990, O.Penrose 1990.

•
$$w = \xi^2, v = u + \frac{l}{2}\varphi \equiv$$
enthalpy function and $\xi \equiv interface width.$

♣ Previous Results. Asymptotic behaviour of the solutions (φ^{ξ}, v^{ξ}) of the system (two different phases)

$$\begin{cases} \tau \varphi_t &= \xi^2 \varphi_{xx} - \frac{1}{2} (\varphi^3 - \varphi) - l\varphi + 2v, \ x \in (a, b) \\ v_t &= k v_{xx} - \frac{kl}{2} \varphi_{xx}, \\ \varphi'(a) &= \varphi'(b) = 0 \\ v'(a) &= v'(b) = 0 \\ \varphi^{\xi}(0, x) &= \varphi_0^{\xi}(x) \in H^1(a, b) \\ v^{\xi}(0, x) &= v_0^{\xi}(x) \in L^2(a, b) \end{cases}$$
(0.6)

when $\xi \sim 0$.

- Metastable Solutions (Nor equilibrium points. Nor energy minima. But
- Have a Slow Evolution(using Energy methods (Cahn-Hilliard, Cahn-Morral system [4, 16]) (Jimenez-Casas[18, 20], Jimenez-Casas-Rodriguez-Bernal [19])
- $u(t, x) \equiv$ the concentration of the point x at time t, of one the components of the mixture.
- We consider the dynamics of phase separation and coarsening of mixtures of three or more components.

 \clubsuit Asymptotic behaviour of the solutions (φ^{ξ},v^{ξ}) of the system

$$\begin{cases} \tau \varphi_t &= \xi^2 \varphi_{xx} - G'(\varphi) - l\varphi + 2v, \ x \in (a, b) \\ v_t &= k v_{xx} - \frac{kl}{2} \varphi_{xx}, \qquad x \in (a, b) \\ \varphi'(a) &= \varphi'(b) = 0 \\ v'(a) &= v'(b) = 0 \\ \varphi^{\xi}(0, x) &= \varphi^{\xi}_0(x) \in H^1(a, b) \\ v^{\xi}(0, x) &= v^{\xi}_0(x) \in L^2(a, b) \end{cases}$$
(0.7)

when $\xi \sim 0$.

• $G'(\varphi) \equiv$ general density function, instead $\frac{1}{2}(\varphi^3 - \varphi)$

•
$$G \ge 0$$
 with $G \in \mathcal{C}^3$

- G has only finitely many zeros, $G^{-1}(0) = \{z_1, ..., z_m\}$ (corresponding to the states or phases of the system).
- $G''(z_i) > 0, i = 1, ..., m$ (in this points G take the minimum.)
- for initial data $(\varphi_0^{\xi}, v_0^{\xi})$, where $\varphi_0^{\xi} \sim z_i$ except at the transition points, and $v_0^{\xi} \sim \frac{l}{2} \varphi_0^{\xi}$.
- Metastable Solutions (Nor equilibrium points. Nor energy minima)
- Have a Slow Evolution

h The Normalized Energy.

Lema 0.1. The energy functional defined by

$$F_{\xi}(\varphi, v) = \int_{a}^{b} \left[\frac{\xi^{2}}{2}\varphi_{x}^{2} + G(\varphi)\right] dx + \frac{l}{2} \int_{a}^{b} \left(\frac{2}{l}v - \varphi\right)^{2} dx \quad (0.8)$$

is a Lyapunov functional for the system (0.7) in $H^1(a, b) \times L^2(a, b)$.

In particular we have that

$$\frac{d}{dt}F_{\xi}(\varphi^{\xi}, v^{\xi}) + (\tau \|\varphi_t^{\xi}\|^2 + d\|[(-\Delta)^{-1}v_t^{\xi}]\|^2) = 0 \qquad (0.9)$$

with
$$d = \frac{4}{kl} > 0$$
.

- $F_{\xi}(\varphi, v) \ge 0.$
- $F_{\xi}(\varphi, v)$ in (0.8) has a shallow valley of energy as $\xi \ll 1$. (Cahn-Hilliard, Cahn-Morral system [4, 16]).
- For initial data in such region little energy is left to be dissipated and thus this translates into a slow evolution in time
- Transitions $(\varphi, v) \equiv \varphi \sim z_i \text{ and } v \sim \pm \frac{l}{2}\varphi$ $\varphi \text{ with large gradients on small transition intervals}$

 $0 \le F_{\xi}(\varphi^{\xi}(t,x), v^{\xi}(t,x)) \le F_{\xi}(\varphi^{\xi}(0,x), v^{\xi}(0,x)) \le h(\xi), \xi << 1$

Definition 1. N, m-transition

A N, m-step with transition points, $y_j, j \in \{1, 2, .., N\}, \varphi^0 : [a, b] \to \{z_i, i = 1, ..., m\},$ $\varphi^0 = \sum_{i=1}^{N+1} z_i \mathcal{X}_{I_i}$ where \mathcal{X} denotes the characteristic function of a set, with

$$I_i \cap I_j = \emptyset$$
, if $i \neq j, \overline{I_1} \cup \overline{I_2} \dots \cup \overline{I_{N+1}} = [a, b]$

$$(\partial(I_1) \cap \partial(I_2) \cap ... \partial(I_{N+1})) \cap (a, b) = \{y_j, j = 1, ..., N\}.$$

(if $N > m - 1, z_{m+r} = z_r, r = 1, 2, N + 1 - m$)

A N, m-transition function is any function in $H^1(a, b)$, which is close to a N, m-step in $L^1(a, b)$.

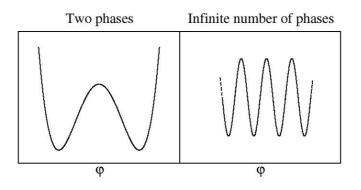


Figure 2: Density function for two phases or m phases

Rescaled Energy Functional

• If
$$\liminf_{\xi \to 0} F_{\xi}(\varphi_0^{\xi}, v_0^{\xi}) \equiv O(\xi^2)$$
, then
 $\varphi_0^{\xi} \equiv z_i \ \acute{o} -1 \quad and \quad v_0^{\xi} \equiv \pm \frac{l}{2} z_i$.

- If we used O(ξ) ≡ instead, we can include a large class of functions (φ^ξ₀, v^ξ₀) [4],[16].
- $V_{\xi} = \frac{1}{\xi} F_{\xi} \equiv$ Rescaled Energy Functional

$$V_{\xi}(\varphi, v) = E_{\xi}(\varphi) + \frac{l}{2\xi} \int_{a}^{b} (\frac{2}{l}v - \varphi)^{2} dx$$

with

$$E_{\xi}(\varphi) = \int_{a}^{b} \left[\frac{\xi}{2}\varphi_{x}^{2} + \frac{1}{\xi}G(\varphi)\right]dx. \qquad (0.10)$$

Lema 0.2. If $\{\varphi^{\xi}\} \subset H^1(a, b)$, such that $\varphi^{\xi} \longrightarrow \varphi^0$ in $L^1(a, b)$ when $\xi \to 0$, and φ^0 a function N, m-step, then:

$$\liminf_{\xi \to 0^+} E_{\xi}[\varphi^{\xi}] \ge \frac{1}{2} \sum_{i=1}^N H^*(z_i + 1) - H^*(z_i) = C(N, m)$$

with $H^*(s) = \int_0^s H(r) dr$ and $H(s) = |2G(s)|^{\frac{1}{2}}$.

3. SLOW MOTION FOR MORE OF TWO DIFFERENT PHASES

- $\varphi^0 \equiv N, m-step function.$
 - $y_j, j = 1, .., N$ are the transition points
 - r is such that $(y_j r, y_j + r) \subset (a, b)$ are disjoint, with $0 < C \leq r$.
- Initial data $\equiv N, m$ -transition ([16]).
- We show an estimate on the norm of this solution in the product space $L^2(a,b) \times H^{-1}(a,b)$.

Proposition 1. We assume that the initial data $(\varphi_0^{\xi}(x), v_0^{\xi}(x))$ is close to the structure of N, m-transition, i.e.

i) $\lim_{\xi \to 0} \varphi_0^{\xi}(x) = \varphi^0(x)$ in $L^1(\Omega).(\varphi^0 \text{ is } N, m-\text{step function})$

ii)
$$E_{\xi}[\varphi_0^{\xi}] \leq C(N,m) + \frac{1}{2}h(\xi)$$
, with $\xi h(\xi) \to 0$ as $\xi \to 0$.
iii) $l \int_a^b |\frac{2}{l}v_0^{\xi} - \varphi_0^{\xi}|^2 dx \leq \xi h(\xi)$.

Then, there exits C_1, C_2 positive constants independent of of ξ , such that the solution (φ^{ξ}, v^{ξ}) satisfies

$$\int_0^T \int_a^b [(\varphi_t^{\xi})^2 + |(-\Delta)^{-1}(v_t^{\xi})|^2] dx dt \le C_1(\xi h(\xi) + \xi e^{-\frac{C}{\xi}})$$

for ξ sufficiently small, and we can choose T such that

$$T \ge \frac{C_2}{C_1(\xi h(\xi) + \xi e^{-\frac{C}{\xi}})}.$$

In particular, if $h(\xi) = C_3 e^{-\frac{C}{\xi}}$, then

$$T \ge C_4 e^{\frac{C}{\xi}}, C_i > 0, i = 3, 4.$$

Slow motion when τ is independent of interface width ξ .

- We assume that the initial data $(\varphi_0^{\xi}(x), v_0^{\xi}(x))$ is close to the structure of N, m-transition.
- The initial structure of N, m-transition solution, is preserved, for a time scale of length T with T ≥ Me^{C/ξ}.
 Teorema 0.3. We assume that the initial data (φ^ξ₀(x), v^ξ₀(x))) satisfies the hypotheses in Proposition 1, i.e.
 i) lim_{ξ→0} φ^ξ₀(x) = φ⁰(x) in L¹(Ω). (φ⁰ is N, m-step function)
 ii) E_ξ[φ^ξ₀] ≤ C(N, m) + ½h(ξ), with ξh(ξ) → 0 as ξ → 0.
 iii) l ∫^b_a |²_lv^ξ₀ φ^ξ₀|²dx ≤ ξh(ξ).
 Then, for any M > 0

$$\begin{split} i) & \lim_{\xi \to 0} \sup_{\{0 \le t \le \frac{M}{h(\xi) + e^{-\frac{C}{\xi}}\}} } \|\varphi^{\xi}(t) - \varphi^{0}\|_{L^{1}} = 0. \\ ii) & \lim_{\xi \to 0} \sup_{\{0 \le t \le \frac{M}{h(\xi) + e^{-\frac{C}{\xi}}\}} } \|\frac{2}{l}v^{\xi}(t) - \varphi^{\xi}(t)\|_{L^{2}} = 0. \\ iii) & \lim_{\xi \to 0} \sup_{\{0 \le t \le \frac{M}{h(\xi) + e^{-\frac{C}{\xi}}\}} } \|\frac{2}{l}v^{\xi}(t) - \varphi^{0}\|_{L^{1}} = 0. \\ In \ particular, \ if \ h(\xi) = ke^{-\frac{C}{\xi}} \ for \ some \ k, \ then \\ iv) & \lim_{\xi \to 0} \sup_{0 \le t \le Me^{\frac{C}{\xi}} } \|\varphi^{\xi}(t) - \varphi^{0}\|_{L^{1}} = 0. \\ v) & \lim_{\xi \to 0} \sup_{0 \le t \le Me^{\frac{C}{\xi}} } \|\frac{2}{l}v^{\xi}(t) - \varphi^{\xi}(t)\|_{L^{2}} = 0. \\ vi) & \lim_{\xi \to 0} \sup_{0 \le t \le Me^{\frac{C}{\xi}} } \|\frac{2}{l}v^{\xi}(t) - \varphi^{0}\|_{L^{1}} = 0. \end{split}$$

♣ Metastable solutions for the thin-interface limit

- Now we study the thin-interface limit, this is, we consider now $\tau = \xi^2$ together with $w = \xi^2$, where ξ (interface width) goes to zero.
- In this case, we consider the initial datum φ₀ very closed to the N, m transition structure. This is, we assume that E_ξ[φ^ξ₀] ≤ C(N, m) + ½h(ξ), with h(ξ) such that ξ⁻¹h(ξ) → 0 as ξ → 0. (instead ξh(ξ) → 0 as ξ → 0 with τ independent of ξ.)
- We prove that the initial structure of N, m transition solution, is preserved for a time scale of length T with

- $T \ge M\xi^{1+\delta}e^{C/\xi}$, for any positive constants M, δ , (instead $T \ge Me^{\frac{C}{\xi}}$).
- Thus, in this case we prove the solutions is preserves during an interval of time smaller than the above case.

Teorema 0.4. We assume that the initial data $(\varphi_0^{\xi}(x), v_0^{\xi}(x))$ satisfy:

$$i) \lim_{\xi \to 0} \varphi_0^{\xi}(x) = \varphi^0(x) \text{ in } L^1(\Omega).$$

$$ii) E_{\xi}[\varphi_0^{\xi}] \leq C(N,m) + \frac{1}{2}h(\xi), \text{ with } \xi^{-1}h(\xi) \to 0 \text{ as } \xi \to 0.$$

$$iii) l \int_a^b |\frac{2}{l} v_0^{\xi} - \frac{1}{2}h(\varphi_0^{\xi})|^2 dx \leq \xi h(\xi).$$

$$Then \text{ for any } M > 0, \delta > 0 \text{ we have}$$

$$i) \lim_{\xi \to 0} \sup_{\{0 \leq t \leq \frac{M\xi^{1+\delta}}{h(\xi)+e^{-\frac{C}{\xi}}\}}} \|\varphi^{\xi}(t) - \varphi^0\|_{L^1} = 0.$$

$$ii) \lim_{\xi \to 0} \sup_{\{0 \leq t \leq \frac{M\xi^{1+\delta}}{h(\xi)+e^{-\frac{C}{\xi}}\}}} \|\frac{2}{l}v^{\xi}(t) - \varphi^{\xi}(t)\|_{L^2} = 0.$$

$$iii) \lim_{\xi \to 0} \sup_{\{0 \leq t \leq \frac{M\xi^{1+\delta}}{h(\xi)+e^{-\frac{C}{\xi}}\}}} \|\frac{2}{l}v^{\xi}(t) - \varphi^0\|_{L^1} = 0.$$

In particular, if
$$h(\xi) = ke^{-\frac{C}{\xi}}$$
 for some k , then
iv) $\lim_{\xi \to 0} \sup_{0 \le t \le M\xi^{1+\delta_e} \frac{C}{\xi}} \|\varphi^{\xi}(t) - \varphi^{0}\|_{L^1} = 0.$
v) $\lim_{\xi \to 0} \sup_{0 \le t \le M\xi^{1+\delta_e} \frac{C}{\xi}} \|\frac{2}{l}v^{\xi}(t) - \varphi^{\xi}(t)\|_{L^2} = 0.$
vi) $\lim_{\xi \to 0} \sup_{0 \le t \le M\xi^{1+\delta_e} \frac{C}{\xi}} \|\frac{2c}{l}v^{\xi}(t) - \varphi^{0}\|_{L^1} = 0.$

4. METASTABLE SOLUTIONS FOR NONLINEAR DIFFUSION PROB-LEM

 \clubsuit Asymptotic behaviour of the solutions (φ^{ξ},v^{ξ}) of the system

$$\begin{cases} \tau \varphi_t &= \xi^p (|\varphi_x|^{p-2} \varphi_x)_x - G'(\varphi) - l\varphi + 2v, \ x \in (a,b) \\ v_t &= k v_{xx} - \frac{kl}{2} \varphi_{xx}, \\ \varphi'(a) &= \varphi'(b) = 0 \\ v'(a) &= v'(b) = 0 \\ \varphi^{\xi}(0,x) &= \varphi_0^{\xi}(x) \in W^{1,p}(a,b) \\ v^{\xi}(0,x) &= v_0^{\xi}(x) \in L^2(a,b) \end{cases}$$
(0.11)

p > 2, when $\xi \sim 0$.

Teorema 0.5. We assume that the initial data
$$(\varphi_0^{\xi}(x), v_0^{\xi}(x))$$

satisfies the hypotheses in Proposition 1, i.e.
i) $\lim_{\xi \to 0} \varphi_0^{\xi}(x) = \varphi^0(x)$ in $L^p(\Omega)$. $(\varphi^0 \text{ is } N, m-\text{step func-tion})$
ii) $E_{\xi}[\varphi_0^{\xi}] \leq C(N, m, p) + \frac{1}{2}h(\xi)$, with $\xi^{p-1}h(\xi) \to 0$ as $\xi \to 0$.
iii) $l \int_a^b |\frac{2}{l}v_0^{\xi} - \varphi_0^{\xi}|^2 dx \leq \xi^{p-1}h(\xi)$.
Then, for any $M > 0$
i) $\lim_{\xi \to 0} \sup_{\{0 \leq t \leq \frac{M\xi^{1-p}}{h(\xi)+e^{-\frac{C}{\xi}}\}}} \|\varphi^{\xi}(t) - \varphi^0\|_{L^1} = 0.$
ii) $\lim_{\xi \to 0} \sup_{\{0 \leq t \leq \frac{M\xi^{1-p}}{h(\xi)+e^{-\frac{C}{\xi}}\}}} \|\frac{2}{l}v^{\xi}(t) - \varphi^0\|_{L^1} = 0.$

In particular, if
$$h(\xi) = ke^{-\frac{C}{\xi}}$$
 for some k , then
iv) $\lim_{\xi \to 0} \sup_{0 \le t \le M\xi^{1-p_e} \frac{C}{\xi}} \|\varphi^{\xi}(t) - \varphi^{0}\|_{L^1} = 0.$
v) $\lim_{\xi \to 0} \sup_{0 \le t \le M\xi^{1-p_e} \frac{C}{\xi}} \|\frac{2}{l}v^{\xi}(t) - \varphi^{\xi}(t)\|_{L^2} = 0.$
vi) $\lim_{\xi \to 0} \sup_{0 \le t \le M\xi^{1-p_e} \frac{C}{\xi}} \|\frac{2}{l}v^{\xi}(t) - \varphi^{0}\|_{L^1} = 0.$
with $E_{\xi}(\varphi) = \int_a^b (\frac{\xi}{p}|\varphi_x|^p + \frac{1}{\xi^{p-1}}G(\varphi))dx.$

• $F_{\xi}(\varphi, v) = \int_{a}^{b} \left(\frac{\xi^{p}}{p} |\varphi_{x}|^{p} + G(\varphi)\right) dx + \frac{l}{2} \int_{a}^{b} \left(\frac{2}{l}v - \varphi\right)^{2} dx$ Lyapunov functional for the system (0.11) in $W^{1,p}(a, b) \times L^{2}(a, b)$.

•
$$V_{\xi}(\varphi, v) = \frac{1}{\xi^{p-1}} F_{\xi}(\varphi, v) = E_{\xi}(\varphi) + \frac{l}{2\xi^{p-1}} \int_{a}^{b} (\frac{2}{l}v - \varphi)^2 dx$$

\$Numerical experiments

In this section we solve the phase-field equations using the Runge-Kutta

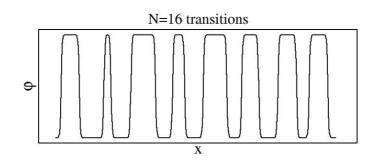
Evolution of phase-field for two phases (m = 2)
 We consider two phases associated to the values +1 and -1.

With this experiments we note that if we consider the initial conditions for φ taking two values +1 and -1 with N = 4 transitions points, this initial structure is conserved for a large interval of time.

We note also the length of interval of time is decreasing when the number of transitions points N is creasing.

This is if we consider $N \geq 4$ then the slow-motion of this

initial structure structure is less than N = 4.



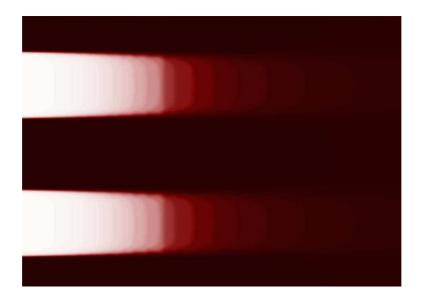


Figure 3: Evolution of phase field for two phases, (time, Phase-field)

• Evolution of phase-field for more than two phases m = 7

In this case we consider m = 7

We note that the solutions with this structure has a slowmotion, this is this initial structure is conserved for a large interval of time.

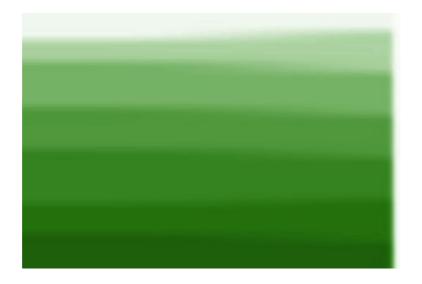


Figure 4: Evolution of phase field for more than two phases, (time, Phase-field)

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